

where  $\alpha$  is defined as the mass fraction of the second phase and  $v_i$  is specific volume of the  $i$ -th phase at  $p$  and  $T$ .

Conditions 4 and 5 imply that the total internal energy of a mass element is the sum of the internal energies of the two phases:

$$E = (1-\alpha)E_1 + \alpha E_2 \quad (3.7)$$

where  $E_i$  is specific internal energy of the  $i$ -th phase at  $p$  and  $T$ .

In order to obtain a suitable constitutive relation when the phase transition has a finite reaction rate, we assume first that the  $p, v, E$  surfaces in each phase can be extended smoothly into metastable regions overlapping the equilibrium region of mixed phase, as in Fig. 3.1.

Then we have the following functional forms:

$$\begin{aligned} v_1 &= v_1(p, T) \\ v_2 &= v_2(p, T) \\ E_1 &= E_1(p, T) \\ E_2 &= E_2(p, T). \end{aligned} \quad (3.8)$$

Then from Eqs. (3.6) and (3.7) we have

$$dv = (1-\alpha)dv_1 + \alpha dv_2 + (v_2 - v_1)d\alpha \quad (3.9)$$

$$dE = (1-\alpha)dE_1 + \alpha dE_2 + (E_2 - E_1)d\alpha. \quad (3.10)$$

From Eq. (3.8):

$$dv_i = \left(\frac{\partial v_i}{\partial T}\right)_p dT + \left(\frac{\partial v_i}{\partial p}\right)_T dp \quad (3.11)$$

$$i = 1 \text{ or } 2.$$

With  $E$  a function of  $T$  and  $p$  we have:

$$\begin{aligned} dE_i &= (\partial E_i / \partial T)_p dT + (\partial E_i / \partial p)_T dp \\ &= (C_{pi} - p(\partial v_i / \partial T)_p) dT - (T(\partial v_i / \partial T)_p + p(\partial v_i / \partial p)_T) dp \end{aligned} \quad (3.12)$$

$$i = 1 \text{ or } 2$$

where  $C_{pi}$  = specific heat of  $i$ -th phase at constant pressure.

Substituting (3.11) and (3.12) into (3.9) and (3.10), we get

$$dv = l_1 dp + m_1 dT + n_1 d\alpha \quad (3.13)$$

$$dE = l_2 dp + m_2 dT + n_2 d\alpha \quad (3.14)$$

where

$$l_1 = (1-\alpha) (\partial v_1 / \partial p)_T + \alpha (\partial v_2 / \partial p)_T \quad (3.15)$$

$$m_1 = (1-\alpha) (\partial v_1 / \partial T)_p + \alpha (\partial v_2 / \partial T)_p \quad (3.16)$$

$$n_1 = v_2 - v_1 \quad (3.17)$$

$$\begin{aligned} l_2 &= -(1-\alpha) (T(\partial v_1 / \partial T)_p + p(\partial v_1 / \partial p)_T) \\ &\quad - \alpha (T(\partial v_2 / \partial T)_p + p(\partial v_2 / \partial p)_T) \end{aligned} \quad (3.18)$$

$$m_2 = (1-\alpha) (C_{p1} - p(\partial v_1 / \partial T)_p) + \alpha (C_{p2} - p(\partial v_2 / \partial T)_p) \quad (3.19)$$

$$n_2 = E_2 - E_1 \quad (3.20)$$

and

$$dE = -(p+q)dv \text{ from Eq. (3.5).} \quad (3.21)$$

Therefore, in principle, (3.13) and (3.14) can be solved for  $dp$  and  $dT$  if  $d\alpha$  is given.

Now we assume the following relaxation